

COLOR DEGREES OF FREEDOM IN NUCLEAR COLLISIONS AND CHARMONIUM SUPPRESSION

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In high energy nuclear collisions, the Glauber model is commonly used to evaluate J/ψ suppression due to inelastic interaction with colorless bound nucleons. This requires an effective value for the J/ψ -nucleon absorption cross section which is larger than theoretically expected. On the other hand, multiple nucleon-nucleon collisions mediated by color exchange interactions, excite their color degrees of freedom. We investigate the importance of this effect and conclude that the related corrections are important to explain the effective value extrapolated from experiment.

In a nucleus-nucleus (AB) collision J/ψ suppression is mostly treated within the Glauber model, with a constant effective cross section for absorption on bound, non-interacting nucleons. Here we intend to study and to improve the understanding of the physical phenomena behind early stage absorption. Following a known interpretation of the dynamics of nucleon-nucleon (NN) collisions, we ascribe the main contribution to the inelastic cross section to color exchange processes. Via color exchange, the 3-quark systems constituting the colliding nucleons become color octet states. Color exchange may also be accompanied by gluon radiation. While the contribution of radiated gluons to J/ψ absorption was treated in ¹, here we focus on the first part. Having realized that color-singlet nucleons are found in color octet states after an elementary collision, one can expect that repeated scatterings lead to multiple color exchanges. One therefore has the possibility to excite the color degrees of freedom of the 3-quark system also to the decuplet state. Higher color representations can be obtained when radiated gluons are also present, but here we restrict our study on the color dynamics of the constituent quarks of a nucleon.

To calculate the value of the inelastic cross section for Ψ on a nucleon which has undergone multiple scattering, we use the two-gluon exchange model of Low and Nussinov ². Although oversimplified, the model incorporates important properties of high-energy hadronic cross sections such as (approximate) energy independence and scaling of the cross sections according to the mean square radius of hadrons. Indicating with Ψ the effective charmonium state, composed by a mixture of J/ψ , ψ' and χ , one can write at once the absorption cross section, obtaining

$$\sigma_{\Psi N_n}^{abs} = \int d^2\vec{k} \frac{1}{(k^2 + \mu^2)^2} \langle \Psi | 1 - \exp(i \vec{k} \vec{x}_{12}) | \Psi \rangle \langle N_n | \Gamma^a(k) \Gamma^{\dagger a}(k) | N_n \rangle. \quad (1)$$

The first factor in the integrand is the squared gluon propagator, the second contains the phase shift due to the momentum transfer \vec{k} and is responsible, by the way, for color transparency, while the last factor is the interesting part of the present calculation. It is the product of two vertices between three quarks and a gluon, averaged with the wave function of the colored nucleon. The vertex is $\Gamma^a(k) = 16/3 \alpha_s [\lambda_{(1)}^a \exp(i \vec{k} \vec{y}_1) + \lambda_{(2)}^a \exp(i \vec{k} \vec{y}_2) + \lambda_{(3)}^a \exp(i \vec{k} \vec{y}_3)]$. In order to average the product of the two vertices, one needs to specify the structure of the

nucleon wave function. The color part of the wave function is the relevant one, while the spatial part is chosen to be equal for all color multiplets. The factor given by the angular brackets in eq. (1) can be evaluated and leads to

$$\langle N_n | \Gamma^a \Gamma^{\dagger a} | N_n \rangle = \langle N | \sum_{i < j} 1 - c_n \exp(i \vec{k} \vec{y}_{ij}) | N \rangle. \quad (2)$$

The effect of having different color states is all contained in the coefficients c_n , listed in the second column of Tab. 1. The values of c_n found, for $n = \mathbf{8}, \mathbf{10}$, lead to larger absorption cross sections as compared to the singlet case. To see this it is useful to define the two-quark form factor $F_N^{2q}(k^2) = \langle N | \exp(i \vec{k} \vec{x}_{ij}) | N \rangle$ for the nucleon. Since, within our approximation, the latter does not depend on the colour multiplet, we drop the label n . Then we can rewrite eq. (1) as $\sigma_{\Psi N_n}^{abs} = \langle \Psi | \sigma_n(r) | \Psi \rangle$, where

$$\sigma_n(r) = \frac{16}{3} \alpha_s^2 \int d^2 \vec{k} \frac{1}{(k^2 + \mu^2)^2} [1 - \exp(i \vec{k} \vec{r})] [1 - c_n F_N^{2q}(k^2)], \quad (3)$$

exhibiting a similar structure to that of the usual dipole-nucleon cross section, but with the important difference carried by the coefficients c_n . Notice that with $c_n \neq 1$ and $\mu = 0$ the integral in eq.(3) is infrared divergent. Therefore, the cut-off μ , of the order of the inverse confinement radius, is important. With $\mu \neq 0$ it holds that $\sigma_n(r) \rightarrow 0$ for $r \rightarrow 0$, so that one still preserves color transparency, even with a colored nucleon. This is indeed a property of the singlet meson. Here we are not primarily interested in the absolute values of $\sigma_{\Psi N_n}^{abs}$ but rather on its dependence on the color state n with respect to the singlet value. Yet we estimated the absolute values choosing $\alpha_s = 0.6$ and $\mu = 140$ MeV. Within a harmonic oscillator model of the nucleon, one finds that $F_N^{2q}(k^2) = F_N^{em}(3k^2)$, the electromagnetic form factor. We assume that this property is general and take the dipole form $F_N^{em}(Q^2) \sim (\lambda^2 + Q^2)^{-2}$ with the parameter $\lambda^2 = 12/R_p^2$ and $R_p = 0.8$ fm. With the above definitions, the generalized dipole-nucleon cross section can be computed analytically and the results of the calculation are illustrated in Fig. ??, showing the modified dipole-nucleon cross section for the three color states of the nucleon. It largely increases when the dipole scatters off a colored object. One can get more insight in the effect of the different color states rewriting eq. (3) in the form

$$\sigma_n(r) = \sigma(r) [1 + (1 - c_n) \Delta(r)] , \quad (4)$$

where $\sigma(r)$ is the usual dipole cross section on a color singlet nucleon and the function $\Delta(r)$ is found to vary appreciably with r . Its functional dependence is given in the small box in Fig. ??. To calculate the inelastic cross section, we must average σ_n with the meson wave function. Therefore, the dependence of Δ on r shows that the effect of having colored nucleons is stronger for larger size mesons. Recent calculations³, based on realistic phenomenology for the dipole cross section and charmonium wave function, gave values for cross sections of various charmonia on a nucleon. At center of mass energy $\sqrt{s} = 10$ GeV, it was found that $\sigma_{J/\psi N} = 3.6$ mb, $\sigma_{\psi' N} = 12.2$ mb and $\sigma_{\chi N} = 9.1$ mb. With a composition of Ψ of 52-60% due to J/ψ , 8-10% due to ψ' and 32-40% due to χ , a weighted value of 5.8 mb was obtained. In the present case we evaluate $\sigma_{\Psi N_n}^{abs}$ with a Gaussian wave function having root mean square transverse separation

Table 1: Results of the calculation of the meson-colored nucleon cross sections.

Multiplet	{n}	c_n	$\sigma_{\psi N_n}$ [mb]	δ_n
Singlet	{ 1 }	1	5.8	1.0
Octet	{ 8 }	-1/4	7.8	1.35
Decuplet	{ 10 }	1/2	9.9	1.7

$\langle r^2 \rangle^{1/2} = 0.25$ fm, in order to account for the different states implicitly included in Ψ , and therefore adjust the singlet cross section to the quoted 5.8 mb. The values found are given in the third column of Tab. 1. In order to minimize the model dependence, we factored out from the colored cross sections the amount corresponding to the singlet case, therefore writing $\sigma_{\psi N_n}^{abs} = \sigma_{\psi N}^{abs} \delta_n$. The values for δ_n are given in the last column of Tab. 1. They constitute the important result of this calculation, showing that Ψ scatters with a color octet nucleon with a 35 % larger cross section with respect to a singlet nucleon, while with the decuplet the increase is even of 70 %. The large values found now require a careful analysis to establish the validity of this newly proposed absorption mechanism.

Since the absorption pattern observed in pA collisions is usually considered to be the baseline for the AB case, it is necessary to examine it in some detail. Before we analyze the data for charmonium production in pA collisions, we emphasize the importance of charmonium formation time⁷. Within a simplified version of a two-channel model⁴, the ΨN absorption cross section evolves in the eigentime τ of Ψ as $\sigma_{\Psi N}^{abs}(\tau) = \sigma_{\Psi N}^{in} + (\sigma^{(0)} - \sigma_{\Psi N}^{in}) \cos(\tau/\tau_f)$, where $\tau_f = (m_{\psi'} - m_{J\psi})^{-1} = 0.3$ fm is the charmonium formation time, $\sigma_{\Psi N}^{in}$ is the inelastic ΨN cross section and $\sigma^{(0)}$ is the cross section of the so called “pre-meson” on a nucleon. As mentioned in the last section, calculations give $\sigma_{\Psi N}^{in} = 5.8$ mb³. Unfortunately $\sigma^{(0)}$ has not yet been calculated. We therefore use the value 2.7 mb from⁴.

One can then write the expression for Ψ production in pA collisions, as

$$\sigma_{\Psi}^{pA} = \sigma_{\Psi}^{NN}/A \int d^2\vec{b} dz \times \left[1 - \int_z^{+\infty} dz' \sigma_{\Psi N}^{abs}(\frac{z'-z}{\gamma_{\Psi}}) \rho_A(z, \vec{b})/A \right]^{A-1} \quad (5)$$

We want to compare this expression with the conventional one obtained in the Glauber model, containing only an effective absorption cross section, therefore making the replacement $\sigma_{\Psi N}^{abs}(\tau) \rightarrow \sigma_{\Psi N}^{eff}$. For $A = 208$ and $p_{\Psi} = 50$ GeV, using a realistic profile for Pb we obtain $\sigma_{\Psi N}^{eff} = 3.8$ mb. This number has to be interpreted as the value of the effective absorption cross section of Ψ on a color singlet nucleon, including formation time effects. To obtain the corresponding value for the case when Ψ scatters on a color octet nucleon we have to enhance the singlet result by 35%, obtaining 5.1 mb.

We have then re-analyzed the existing data taken by 3 fixed target experiments: NA3 at CERN⁵, NA38/51⁶ at CERN and E866 at FERMILAB⁷. In order to remove the effects on absorption due to different Ψ momenta, we impose the same kinematic conditions, considering Ψ 's of 50 GeV momentum in the laboratory frame for the three experiments. We then assume that the absorption mechanism is parameterized by the effective cross section $\sigma_{\Psi N}^{eff}$. We made three separate fits including quoted statistical and systematic errors and extracted the effective absorption cross section by minimizing χ^2 . The result of these fits is presented in the second column of Tab. 2. The mean value of 3.5 ± 0.6 mb seems to favor the singlet mechanism, for which a cross section of 3.8 mb is theoretically expected.

Table 2: Results of the fits to pA data for the various experiments considered. See text for details.

Experiments $p_{\psi} = 50$ GeV	$\sigma_{\psi N}^{abs}$ [mb]	
	Glauber fit	Gluon corr.
NA3	4.5 ± 1.7	7.3 ± 1.9
NA38/51	7.1 ± 1.6	10 ± 1.8
E866	2.7 ± 0.7	4.7 ± 0.7
Averages	3.5 ± 0.6	5.7 ± 0.6
Theory	$3.8 \{1\}$	$5.1 \{8\}$

On the other hand, the assumption that Ψ is produced on a bound nucleon with the same cross section as on a free one, can be an oversimplification of the physics under study. In a nuclear environment parton distributions are substantially modified. In particular, the gluon distribution in nuclei exhibits anti-shadowing effects at $x_{Bj} \simeq 0.1$. We expect that charmonium production on a nucleus is enhanced by a factor $R_A^g(x_2, Q^2) = g_A(x_2, Q^2)/g(x_2, Q^2)$. We estimate these effects by rescaling the Ψ cross section with the A-dependent correction $R_A^g = 0.032 \log(A) + 1.006$, extracted from⁸. We fitted the data once more, obtaining the results listed in the third column of Tab. 2. One sees that the new average value of 5.7 ± 0.6 mb now seems to favor the octet mechanism, for which we estimated a value of 5.1 mb. One sees that the situation is, in all, quite uncertain.

The above analysis assumes that there is no dependence of nuclear effects on the proton energy if the charmonium energy is the same in all experiments. However, there might still be some dependence on the proton energy. According to⁹, energy loss of projectile partons in the nuclear medium substantially reduces the production rate of charmonia at large x_F and is still rather important at small x_F . In order to compensate for energy loss, the gluon distribution must be evaluated at a higher value of x_1 , therefore suppressing Ψ production. This will in turn reduce the final state absorption on nucleons in a fit to the data. To account for this effect, further study is required.

To summarize, the data seem to suggest that Ψ is absorbed by colored nucleons in pA collisions. The picture is though not yet conclusive due to the large experimental errors. Nevertheless, we follow our theoretical prediction and examine the potential role of the color mechanism in the more complex case of AB collisions. Preparing the ground for the forthcoming discussion, it is useful to recall the conventional expression within the Glauber model for the Ψ production cross section at impact parameter \vec{b} . One has

$$\begin{aligned} \frac{d^2\sigma_{\Psi}^{AB}}{d^2\vec{b}}(\vec{b}) &= \int d^2\vec{s} dz_A dz_B \sigma_{\Psi}^{NN} \rho_A(z_A, \vec{s}) \rho_B(z_B, \vec{b} - \vec{s}) \\ &\times \left[1 - \sigma_{\Psi N}^{abs} T_A(z_A, \vec{s})/A\right]^{A-1} \left[1 - \sigma_{\Psi N}^{abs} T_B(z_B, \vec{b} - \vec{s})/B\right]^{B-1}. \end{aligned} \quad (6)$$

Here we neglect the effects of any additional suppression factors which go beyond conventional nuclear effects, that is final state interactions with the produced matter (QGP, hadron gas, etc.), since we are not primarily interested in these contributions. On the other hand, we focus on the reinterpretation of the terms describing nuclear effects discussing color exchange in the context of multiple scattering, then applying the obtained results to the calculation of Ψ production in AB collisions.

Having completed the computation of the Ψ cross sections with colored nucleons, we need to look at how these color objects develop in a collision between nuclei. Consider a nucleon of one of the nuclei colliding with a row of nucleons of the other nucleus. It undergoes multiple scattering, with repeated color exchanges. It is possible to calculate the probability $P_n(z, \vec{b})$ for a nucleon in color state $n = \mathbf{1}, \mathbf{8}, \mathbf{10}$ after having traveled through the row up to z at impact parameter \vec{b} . This can be done in a very elegant way, by solving an evolution equation for the color density matrix of the 3-quark system¹⁰. One obtains

$$P_n(z, \vec{b}) = \frac{1}{27} \times \begin{cases} 1 + 20 F(z, \vec{b}) + 6 G(z, \vec{b}) & \{\mathbf{1}\} \\ 16 - 40 F(z, \vec{b}) + 24 G(z, \vec{b}) & \{\mathbf{8}\} \\ 10 + 20 F(z, \vec{b}) - 30 G(z, \vec{b}) & \{\mathbf{10}\} \end{cases}, \quad (7)$$

where $F(z, \vec{b}) = \exp(-\frac{9}{8} \sigma_{NN}^{in} T(z, \vec{b}))$ and $G(z, \vec{b}) = \exp(-\frac{3}{4} \sigma_{NN}^{in} T(z, \vec{b}))$. One notices several properties of the probabilities. First of all, the limit $T \rightarrow 0$ implies $F, G \rightarrow 1$. This means that $P_1 \rightarrow 1$ and $P_8, P_{10} \rightarrow 0$. In other words the nucleon, before entering the nucleus, is in singlet

state as it should be. Moreover, $P_1 + P_8 + P_{10} = 1$ for any T , therefore probability is always conserved. Finally, if the nucleus is large enough one has $T \gg 1$ and $F, G \ll 1$. This implies that if sufficient scatterings have taken place, the color probabilities reach the statistical limit, given by the first coefficients of eqs. (7). It turns out that soon after the nucleon has penetrated the nuclear profile, the statistical limit is essentially reached in which about $16/27 \simeq 2/3$ of the nucleons are in octet state while $10/27 \simeq 1/3$ are in decuplet. The amount of singlet is negligible.

With the calculated probabilities P_n and the previously evaluated color cross sections $\sigma_{\Psi N_n}^{abs}$, we can now construct the cross section for Ψ when it scatters with a nucleon, taking into account all its color states, and modify eq. (6) by making the replacement

$$\sigma_{\Psi N}^{abs} T_A(z_A, \vec{b}_A) \rightarrow \int_{-\infty}^{z_A} dz'_A \Sigma_B(\tilde{z}_B, \vec{b}_B) \rho_A(z'_A, \vec{b}_A), \quad \text{where} \quad \tilde{z}_B = z_B - (z'_A - z_A) \frac{1 + v_\Psi}{1 - v_\Psi}. \quad (8)$$

The position-dependent effective cross section

$$\Sigma_B(\tilde{z}_B, \vec{b}_B) = \sum_n \sigma_{\Psi N_n}^{abs} P_n^B(\tilde{z}_B, \vec{b}_B), \quad (9)$$

represents absorption by nucleus B , taking into account the multiple scattering of its nucleons with nucleus A . Analogous expressions hold for A .

The modified version of eq. (6) can now be used to evaluate total cross sections, with a further integration in impact parameter. We compute values to compare with available data measured by the NA38/50/51 collaboration at the SPS¹¹. We use realistic nuclear profiles and scale the center of mass energy according to the parameterization $\sigma_{\Psi}^{NN} \sim (1 - m_{J/\psi}/\sqrt{s})^{12}$. The results are compared with the data in Fig. 1. The dashed curve corresponds to a standard Glauber model calculation with the singlet absorption cross section $\sigma_{\Psi N_1}^{abs} = 3.8$ mb, as evaluated in Sect. 3, while the thick curve refers to the improved calculation which takes into account color excitations, studied in Sect. 4. In order to emphasize the difference between the two absorption mechanisms, we neglect here anti-shadowing effects. For the pA calculation we increased the value of $\sigma_{\Psi N_1}^{abs}$ by 35 %, therefore used $\sigma_{\Psi N_8}^{abs} = 5.1$ mb. This is the extreme situation, when Ψ always encounters color octet nucleons. For the AB case we used eq. (9).

Different from the Glauber model result, our improved one provides a better agreement with the data, except for the Pb+Pb point. In other words, the proposed absorption mechanism due to color dynamics is found to be an important contribution to nuclear effects in Ψ production. It must be accounted for, together with other mechanisms of suppression which we here put aside.

To summarize, we studied the dynamics underlying Ψ suppression due to nuclear effects, employing the idea that color exchange processes in multiple NN collisions may play a significant role. We calculated inelastic cross sections for Ψ scattering on colored nucleons and found an increase of about 35 % when the nucleon is in a color octet state as compared with the singlet case, while the enhancement factor for the color decuplet is about 70 %.

Making use of available calculations of ΨN cross sections, we included formation time effects and obtained the effective value for Ψ absorption on singlet nucleons $\sigma_{\Psi N_1}^{abs} = 3.8$ mb. Scaling this value with a 35 % increase we obtained $\sigma_{\Psi N_8}^{abs} = 5.1$ mb and with a 70 % increase $\sigma_{\Psi N_{10}}^{abs} = 6.4$ mb. We then performed a fit to available data on Ψ production cross sections in pA collisions. The result of the fit, including gluon anti-shadowing corrections, is that the effective absorption cross section for Ψ is $\sigma_{\Psi N}^{eff} = 5.7 \pm 0.6$ mb, therefore suggesting that in pA collisions, the produced charmonium interacts with color octet nucleons. The conclusion is, though, rather weak because the experimental points are quite scattered.

Total production cross sections for Ψ were then calculated and compared in Fig. 1 with available data for AB collisions. Compared to the Glauber model result, the agreement is

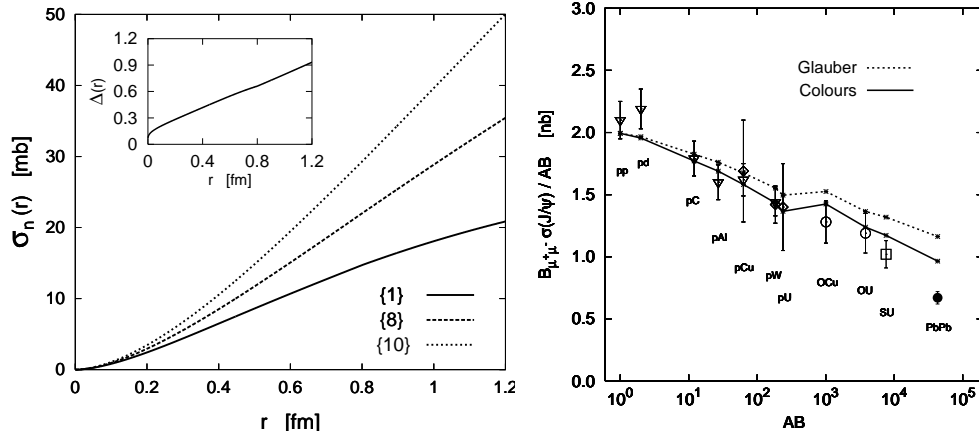


Figure 1: Left: Dipole cross section for different color states as a function of the quark separation r . Box: The quantity $\Delta(r)$ which governs the dependence on color. Right: Calculations of total J/ψ production cross sections in AB collisions and comparison with the NA38/50/51 data. The curves, drawn to guide the eye, correspond to Glauber absorption (dashed) and to absorption by colored nucleons (full).

improved and we conclude that effects due to color exchange processes are substantial and provide the value $\sigma_{\Psi N}^{abs} = 5.5$ mb for the effective absorption cross section on nucleons in AB collisions (See eq. (9)).

In conclusion, we have completed a study of nuclear effects in charmonium production in AB collisions, providing a better baseline for future calculations. These amount to an improved treatment of gluon bremsstrahlung, including non-linearities arising because of gluon cascades, to the calculation of effects caused by gluon energy loss and to the evaluation of E_{\perp} -dependent cross sections. This will hopefully provide the starting point from which to address more precisely the nature of anomalous suppression in Pb+Pb collisions.

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